

# The Vibrational Spectra of Tetrafluorosuccinic Anhydride

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The IR spectra ( $4000\text{--}200\text{ cm}^{-1}$ ) of tetrafluorosuccinic anhydride as a vapour, liquid and dissolved in  $\text{CS}_2$  were recorded. Additional spectra from the vapour were obtained in the region  $400\text{--}30\text{ cm}^{-1}$ . Raman spectra of the liquid were recorded, and semiquantitative polarization ratios were measured. The fundamental frequencies were tentatively assigned in terms of  $C_2$  symmetry based upon Raman polarization data and vapour band contours.

The IR spectra of succinic anhydride as a solid and dissolved in various solvents were recorded between  $4000$  and  $30\text{ cm}^{-1}$ . Raman spectra of the solid were also recorded. The fundamental frequencies were assigned in terms of  $C_{2v}$  symmetry and compared to those of tetrafluorosuccinic anhydride.

Force fields were derived by the overlay technique for succinic and tetrafluorosuccinic anhydride.

We have for some time been interested in the vibrational spectra of five-membered, unsaturated cyclic molecules, and have earlier reported our results for maleimide,<sup>1</sup> *N*-deuterio<sup>1</sup> and *N*-chloro<sup>2</sup> maleimide, succinimide<sup>3</sup> and *N*-deuterio,<sup>3</sup> *N*-chloro, *N*-bromo and *N*-iodo succinimide.<sup>4</sup> These studies have now been extended to include succinic anhydrides. In the present paper we shall report our results for tetrafluorosuccinic anhydride (FSA). Various authors<sup>5–7</sup> have discussed the IR and Raman spectra of succinic anhydrides, but to our knowledge no IR or Raman data have been reported for FSA.

In the present work we have recorded IR spectra of FSA as a vapour, liquid and dissolved in  $\text{CS}_2$ . Raman spectra were obtained from the neat liquid. Unfortunately we were not able to record spectra of FSA in the solid state since the sample seemed to polymerize upon crystallization.

As an aid to the assignments normal coordinate analyses of FSA were carried out. The data were fitted together with the frequencies from the solution state of succinic anhydride (SA). As the earlier IR

data for SA in the solution are rather incomplete, we have also recorded IR spectra of SA dissolved in  $\text{CCl}_4$ ,  $\text{CS}_2$  and  $\text{CH}_2\text{Cl}_2$ .

## EXPERIMENTAL

The sample of FSA was a commercial product from the K & K laboratories. No impurities were detected by gas chromatography. SA was a commercial product from Fluka AG and was purified by repeated sublimation on a cold finger.

Raman spectra of the pure liquid FSA were recorded with a Cary Model 81 spectrometer, modified for perpendicular illumination. The  $5145\text{ \AA}$  line of an argon ion laser was used for excitation. IR spectra were recorded on a Perkin-Elmer Model 225 spectrometer in the region  $5000\text{--}200\text{ cm}^{-1}$ . Far IR spectra were recorded of the vapour (FSA) with an RIIC interferometer FS-520, equipped with a light pipe gas cell of optical path *ca.* 6 m. FIR spectra ( $400\text{--}30\text{ cm}^{-1}$ ) of SA as a polyethylene pellet and dissolved in  $\text{C}_6\text{H}_6$  were recorded using a Perkin-Elmer, Hitachi model FIS-3 spectrometer.

## RESULTS

FSA has recently been studied by electron diffraction.<sup>8</sup> In contrast to SA<sup>9</sup> it was found to have a non-planar ring skeleton although planarity was found for the  $\text{O}=\text{C}-\text{O}-\text{C}=\text{O}$  group. Hence the spectra of FSA will be interpreted in terms of  $C_2$  symmetry, while SA is considered to possess  $C_{2v}$  symmetry.

The 27 fundamentals of SA will accordingly divide themselves into the symmetry species:  $9 a_1 + 8 b_2 + 5 a_2 + 5 b_1$ . The  $a_1$  and  $b_2$  fundamentals represent in-plane, the  $a_2$  and  $b_1$  out-of-plane modes. All the fundamentals are active in the Raman effect while those of species  $a_2$  are forbidden in the infrared spectra.

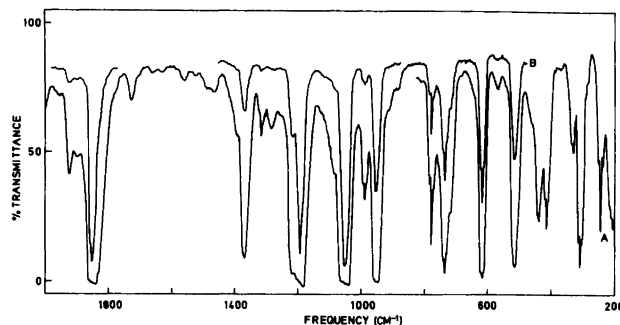


Fig. 1. The infrared spectrum of tetrafluorosuccinic anhydride as a vapour. Path length, 10 cm. Pressure, A, 100 Torr; B, 25 Torr.

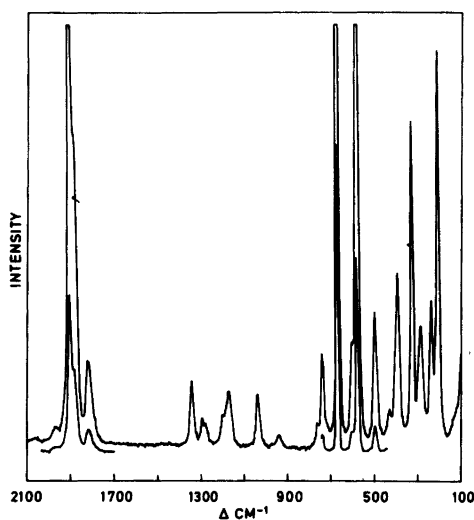


Fig. 2. The Raman spectrum of tetrafluorosuccinic anhydride as a neat liquid.

For FSA all the fundamentals are active both in the Raman and in the IR spectra, and the 27 fundamentals will divide themselves into the symmetry species: 14 *a*+13 *b*. The *a* species correspond to  $a_1 + a_2$  and the *b* species to  $b_1 + b_2$  for SA.

From the structural parameters<sup>8</sup> of FSA the following moments of inertia are obtained ( $10^{-40}$  g cm<sup>2</sup>):  $I_A = 331.7$ ,  $I_B = 412.7$  and  $I_C = 545.8$ . The following PR separations (cm<sup>-1</sup>) were calculated<sup>10</sup> at 300 K: 10 for *A*, 11 for *B* and 14 for *C*-bands. The *Q-Q* separation for the *B*-bands should be 2.5 cm<sup>-1</sup>. Vibrations belonging to species *a* give rise to *B*-bands in the IR vapour spectra, while vibrations of species *b* give *A/C* hybrid bands. Thus, the IR vapour band contours and the Raman polarization measurements offered reasonable good criteria for the assignments of most of the fundamentals for FSA.

*Spectral interpretations.* As an illustration the IR vapour spectrum of FSA is shown in Fig. 1, whereas a Raman spectrum of the liquid appears in Fig. 2.

Table 1. Infrared<sup>a</sup> and Raman spectral data for tetrafluorosuccinic anhydride (FSA).

Infrared		Raman	Assignments
Vapour	Liquid	Liquid	
2005 m <sup>b</sup>	1997 m		$\nu_{14} + \nu_{15}$ b
		1987 vw P	$\nu_{15} + \nu_{27}$ a
1955 w	1955 w sh		$\nu_4 + \nu_{20}$ b
1923 s	1923 m	1925 s P	$\nu_1$ a
1899 m	1903 m	1903 m P	$\nu_{16} + \nu_{22}$ a
	1887 w sh		$\nu_3 + \nu_7$ a
1858 s sh)	1840 vs	1840 m D	$\nu_{15}$ b
1852 s )			
1847 sh )			

Table 1. Continued.

1830 s sh	~ 1815 s sh	1812 vw sh	$\nu_9 + \nu_{16}$	b
1770 w sh			$\nu_2 + \nu_{23}$	b
1732 sh	1740 w sh		$\nu_6 + \nu_{19}$	b
1728 m	1723 m		$\nu_3 + \nu_9$	a
1487 m	1484 w		$\nu_4 + \nu_{24}$	b
1466 m	1462 w		$2 \times \nu_{21}$	a
1395 m sh	1390 m sh		$\nu_7 + \nu_{21}$	b
1369 s	1361 s	1361 m P	$\nu_2$	a
1347 w sh			$\nu_{18} + \nu_{24}$	a
1323 w sh				
1317 m	1313 m	1314 w D	$\nu_{16}$	b
1310 w sh				
1303 w	1300 w		$\nu_4 + \nu_{27}$	b
1293 vw sh				
1285 w	1283 w		$\nu_{12} + \nu_{19}$	b
1277 vw sh				
1250 vw			$\nu_5 + \nu_{26}$	b
1218 s B	1217 s	1217 w P	$\nu_3$	a
1199 sh				
1194 s	~ 1180 vs	1189 m P	$\nu_4, \nu_{17}$	a,b
1187 sh				
1131 vw	1129 w		$\nu_9 + \nu_{22}$	b
1115 vw	1112 w		$\nu_7 + \nu_{10}$	a
1086 m	~ 1085 m sh		$\nu_{20} + \nu_{24}$	a
1079 m				
1053 vs	~ 1052 s sh	1054 m P	$\nu_5$	a
1048 s	1035 s		$\nu_{18}$	b
1044 sh				
993 m				
989 m	991 s		$\nu_6 + \nu_{25}$	b
983 sh				
954 s	~ 960 s sh	957 w D	$\nu_{19}$	b
951 s	950 s		$\nu_9 + \nu_{10}$	a
919 vw sh				
912 w	912 w sh		$\nu_{12} + \nu_{22}$	b
905 vw sh				
884 w sh			$\nu_7 + \nu_{26}$	b
835 vw			$\nu_{10} + \nu_{23}$	b
807 vw			$\nu_7 + \nu_{27}$	b
783 m				
778 s	776 m	777 w D	$\nu_{20}$	b
770 m sh				
742 s sh	~ 745 w sh	747 s P	$\nu_6$	a
735 s	732 s	728 vw sh	$\nu_{21}$	b
728 s sh				
717 m sh			$\nu_8 + \nu_{27}$	b
682 vw	683 w	683 s P	$\nu_7$	a
620 s sh				
616 s	619 s	618 w	$\nu_{22}$	b
611 s sh	593 w sh	593 s P	$\nu_8$	a
567 w			$2 \times \nu_{13}$	a
544 vw sh	544 vw sh		$\nu_{12} + \nu_{25}$	b
521 sh				
518 sh	509 s	509 m P	$\nu_9$	a
514 s				
460 w sh			$\nu_{13} + \nu_{14}$	a

Table 1. Continued.

~442 s sh } B	439 m	440 w P	$\nu_{10}$	a
437 s				
418 m sh } A/C	403 m	403 m D	$\nu_{23}$	b
413 s				
408 m sh }				
385 vw			$\nu_{14} + \nu_{26}$	b
370 vw			$\nu_{25} + \nu_{27}$	a
335 m } B	337 m	337 s P	$\nu_{11}$	a
329 m }				
311.5 s } A/C		305 w sh P	$\nu_{12}$	a
306.5 s }		291 m D	$\nu_{24}$	b
301.5 s }				
290 vw sh } B			$\nu_{13}$	a
283 vw sh }				
249 s } A/C		248 w D	$\nu_{25}$	b
242.5 s }				
236 s }				
214.5 m } A/C		217 s D	$\nu_{26}$	b
204 s }			$\nu_{14}$	a
180 vw				
135 sh } A/C		140 vw	$\nu_{27}$	b
129 m }				
125 sh }				

<sup>a</sup> Weak infrared bands in the regions 4000–2000  $\text{cm}^{-1}$  and 1700–1500  $\text{cm}^{-1}$  are omitted. <sup>b</sup> w, weak; m, medium; s, strong; v, very; sh, shoulder; P, polarized; D, depolarized; A, B and C, vapour band contours.

Table 2. Observed and calculated fundamental frequencies<sup>a</sup> for tetrafluorosuccinic anhydride (FSA).

Species and No.	Obs.	Calc.	PED <sup>b</sup>
a $\nu_1$	1923	1923	82s <sup>c</sup>
$\nu_2$	1369	1362	46d
$\nu_3$	1218	1233	541 <sub>2</sub> + 421 <sub>1</sub> + 18 $\phi_2$
$\nu_4$	1189	1174	65a
$\nu_5$	1053	1061	69r + 36d
$\nu_6$	747	756	321 <sub>2</sub> + 271 <sub>1</sub>
$\nu_7$	682	674	45 $\pi$ + 25 $\tau_2$
$\nu_8$	593	607	23 $\beta$ + 22 $\alpha$
$\nu_9$	514	552	37 $\theta$
$\nu_{10}$	440	445	54 $\tau_2$ + 32 $\psi_2$ + 17 $\psi_1$
$\nu_{11}$	332	345	29 $\eta$ + 24 $\xi$
$\nu_{12}$	305	309	27 $\phi_2$ + 24d + 15 $\theta$
$\nu_{13}$	285	271	90 $\phi_2$ + 79 $\phi_1$
$\nu_{14}$	180	166	48 $\psi_2$ + 38 $\tau_1$ + 36 $\psi_1$
b $\nu_{15}$	1852	1851	85s
$\nu_{16}$	1317	1297	40r + 341 <sub>1</sub> + 16e
$\nu_{17}$	1194	1190	611 <sub>2</sub> + 251 <sub>1</sub>
$\nu_{18}$	1048	1059	71a + 16 $\xi$
$\nu_{19}$	954	959	261 <sub>1</sub> + 251 <sub>2</sub> + 21a
$\nu_{20}$	778	792	27 $\eta$
$\nu_{21}$	735	703	35 $\theta$
$\nu_{22}$	616	622	68 $\pi$
$\nu_{23}$	413	409	19 $\zeta$ + 15r

Table 2. Continued.

$\nu_{24}$	306	327	$21\psi_1 + 18\varepsilon + 15\theta$
$\nu_{25}$	242	238	$43\tau_1 + 27\tau_3 + 18\psi_1 + 18\psi_2$
$\nu_{26}$	204	204	$75\psi_2 + 69\psi_1 + 15\phi_1$
$\nu_{27}$	129	144	$102\phi_2 + 71\phi_1 + 22\tau_3 + 17\varepsilon$

<sup>a</sup> When possible, frequencies are taken from the vapour phase data. <sup>b</sup> The potential energy distribution defined as  $X_{ik} = 100F_{ii}L_{ik}^2/\lambda_k$ . <sup>c</sup> For meaning of symbols, see Fig. 3. Terms below 15 are neglected.

Table 3. Observed and calculated fundamental frequencies<sup>a</sup> for succinic anhydride (SA).

Species and No.	Obs.	Calc.	PED <sup>b</sup>
$a_1$	$\nu_1$	2953	981 <sup>c</sup>
	$\nu_2$	1872	80s
	$\nu_3$	1432	1447
	$\nu_4$	1268	1288
	$\nu_5$	1197	1169
	$\nu_6$	999	991
	$\nu_7$	813	807
	$\nu_8$	620	607
	$\nu_9$	407	395
$a_2$	$\nu_{10}$	2975	2977
	$\nu_{11}$	1225	1239
	$\nu_{12}$	1005	1012
	$\nu_{13}$	652	637
	$\nu_{14}$	279	251
$b_1$	$\nu_{15}$	2985	2971
	$\nu_{16}$	1160	1132
	$\nu_{17}$	813	832
	$\nu_{18}$	528	551
	$\nu_{19}$	135	158
$b_2$	$\nu_{20}$	2945	2960
	$\nu_{21}$	1796	1802
	$\nu_{22}$	1414	1400
	$\nu_{23}$	1292	1289
	$\nu_{24}$	1047	1072
	$\nu_{25}$	906	924
	$\nu_{26}$	645	659
	$\nu_{27}$	559	559

<sup>a</sup> When possible frequencies are taken from solution. <sup>b</sup> The potential energy distribution defined as  $X_{ik} = F_{ii}L_{ik}^2/\lambda_k$ . <sup>c</sup> For meaning of symbols, see Fig. 3.  $l=l_1=l_2$ ,  $\phi=\phi_1=\phi_2$ ,  $\psi=\psi_1=\psi_2$ . Terms below 15 are neglected.

The observed IR and Raman frequencies are listed in Table 1, while the assigned fundamentals are given in Table 2 (FSA) and Table 3 (SA) together with the calculated frequencies and the potential energy distribution (PED) among the valence coordinates.

Our assignments of the fundamentals of the in-plane modes ( $a_1$  and  $b_2$ ) of SA agree with those previously noted.<sup>5</sup> No complete analyses of the out-of-

plane vibrations are reported. We have interpreted the fundamentals of species  $a_2$  and  $b_1$  by comparing the spectra with those of the succinimides.<sup>3,4</sup> The  $\text{CH}_2$  rock fundamental  $\nu_{12}$  of species  $a_2$  is assigned to a weak band at  $1005\text{ cm}^{-1}$  only observed in the IR spectra of the solid SA. Similar weak bands at about the same frequency are also observed for all the succinimides,<sup>3,4</sup> and we suggest that these

bands should be taken as the  $\text{CH}_2$  rock frequency rather than that previously indicated at *ca.*  $920\text{ cm}^{-1}$ .<sup>4</sup>

It is well-known from the maleimides,<sup>1,2</sup> succinimides,<sup>3,4</sup> maleic<sup>11-13</sup> and succinic<sup>5,6</sup> anhydrides that the  $\text{C}=\text{O}$  stretching of species  $a_1$  is invariably at higher frequency than the  $b_2$  mode. The same rule applies to FSA, and  $\nu_1(a)$  is found at  $1923\text{ cm}^{-1}$  *ca.*  $70\text{ cm}^{-1}$  higher than  $\nu_{15}(b)$  at  $1852\text{ cm}^{-1}$ . A similar frequency splitting is also observed for SA where the two fundamentals are found at *ca.*  $50\text{ cm}^{-1}$  lower than for FSA, at  $1872\text{ cm}^{-1}$  ( $a_1$ ) and  $1796\text{ cm}^{-1}$  ( $b_2$ ). As was earlier discussed for the succinimides,<sup>4</sup> the  $\text{C}=\text{O}$  stretching vibration of species  $b_2$  shows significantly larger frequency shifts upon solution than the symmetric vibration ( $a_1$ ). This is also valid for FSA, and the vapour band at  $1852\text{ cm}^{-1}$  is shifted towards  $1840\text{ cm}^{-1}$  in  $\text{CS}_2$  solution, while the corresponding shift for the vapour band at  $1923\text{ cm}^{-1}$  is only  $4\text{ cm}^{-1}$ .

Several very strong bands in the IR spectra of FSA are observed between  $1400$  and  $700\text{ cm}^{-1}$ . Eight fundamentals involving  $\text{C}-\text{C}$ ,  $\text{C}-\text{O}$  and  $\text{C}-\text{F}$  stretching are to be expected in this region. The  $\text{C}-\text{O}$  stretching vibrations are observed at slightly higher frequencies for SA than for FSA. The symmetric mode ( $a$ ) is found at  $1197$  and  $1189\text{ cm}^{-1}$ , and the asymmetric ( $b$ ) at  $1047$  and  $1035\text{ cm}^{-1}$  for SA in solution and liquid FSA, respectively. As indicated by the PED of Tables 2 and 3, the  $\text{C}-\text{O}$  stretchings are well-localized vibrations although for SA they are somewhat mixed with the  $\text{C}-\text{C}$  stretchings.

From Table 2 four fundamentals,  $\nu_3$  ( $1218$ ),  $\nu_6$  ( $747$ ),  $\nu_{17}$  ( $1194$ ) and  $\nu_{19}$  ( $954\text{ cm}^{-1}$ ), can be interpreted as almost pure  $\text{C}-\text{F}$  stretching vibrations with only small mixings with the  $\text{C}-\text{C}$  stretchings. The four vibrations assigned to  $\text{C}-\text{F}$  stretching correspond to either relatively strong bands in the IR spectra of the vapour or to polarized Raman bands. This is in agreement with a non-planar structure of FSA as was also found from the electron diffraction investigation.<sup>8</sup> If the molecule possessed  $C_{2v}$  symmetry, only three of the four  $\text{C}-\text{F}$  stretchings should give IR active vibrations, and only one should give a polarized Raman band.

The ring and  $\text{C}=\text{O}$  in-plane bending vibrations will be expected in the region  $800-400\text{ cm}^{-1}$ . For FSA two additional FCF bending vibrations at  $735$  and  $514\text{ cm}^{-1}$  are complicating this spectral area. For all the similar five-membered rings investigated<sup>1-6,11-13</sup> a characteristic, usually weak

IR and strong Raman band is observed in the narrow frequency region  $600-650\text{ cm}^{-1}$ . This band has been assigned to an in-plane ring bending, mostly involving the  $\text{C}-\text{O}-\text{C}$  part of the molecule. The corresponding bands of FSA and SA ( $\nu_8$ ) are observed at  $593\text{ cm}^{-1}$  and at  $620\text{ cm}^{-1}$ , respectively.

The symmetric  $\text{C}=\text{O}$  in-plane bending vibration was observed at *ca.*  $400-420\text{ cm}^{-1}$  for maleimide,<sup>1</sup> maleic anhydride,<sup>11</sup> succinimide,<sup>3</sup> SA and their deuterated species.<sup>1,3,5,11</sup> For FSA this mode is situated at  $332\text{ cm}^{-1}$ . A considerable frequency shift for this vibration was also noticed for the halogenated maleimides,<sup>2</sup> succinimides<sup>4</sup> and maleic anhydrides.<sup>12,13</sup> This shift can, except for FSA, be explained by a coupling between the  $\text{C}=\text{O}$  in-plane bending and the  $\text{N}-\text{X}$  or  $\text{C}-\text{X}$  stretching vibrations. For FSA a mixing with the FCF bending vibration at  $514\text{ cm}^{-1}$  could explain the low frequency of the symmetric  $\text{C}=\text{O}$  bending vibration although such a mixing is not calculated by the PED (Table 2).

The far IR spectra of FSA are quite different from those of SA due to the many CCF bending vibrations in this region. All bands below  $350\text{ cm}^{-1}$  (FSA) are considered to be fundamentals (Table 1). The IR vapour band contours and the Raman polarization data make the assignments in this spectral area rather straight-forward.

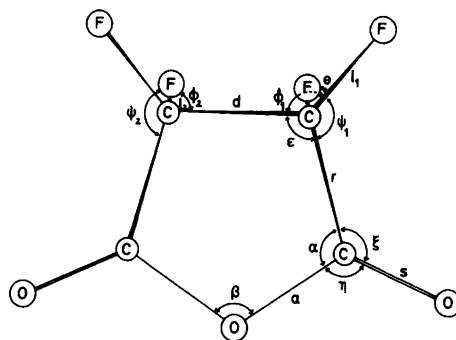


Fig. 3. Internal coordinates for tetrafluorosuccinic anhydride. The following coordinates are not shown on the figure:  $\pi$ , the bending of  $\text{C}'=\text{O}$  relative to the plane formed by the atoms  $\text{C}-\text{C}'-\text{O}$ .  $\tau_1$ ; the normalized sum of the two torsions around the  $\text{C}-\text{O}$  bond.  $\tau_2$ ; the normalized sum of the nine torsions around the  $(\text{F}_2)\text{C}-\text{C}(\text{F}_2)$  bond.  $\tau_3$ ; the normalized sum of the six torsions around the  $(\text{O})\text{C}-\text{C}(\text{F}_2)$  bond.

Table 4. Valence force constants for succinic anhydride (SA) and tetrafluorosuccinic anhydride (FSA).

Force type	Constants		
	Symbol <sup>a</sup>	Value <sup>b</sup>	
		FSA	SA
Stretch	$K_a$	5.68	— <sup>c</sup>
	$K_s$	12.6	11.8
	$K_r = K_d$	5.0	—
	$K_1$	5.5	4.8
Bend	$H_\eta = H_\xi$	0.92	—
	$H_\psi = H_\phi$	1.07	0.57
	$H_\theta$	1.85	0.5
	$H_\alpha = H_\beta = H_\epsilon$	1.33	—
Outo	$\pi$	0.44	—
Torsion	$\tau_1$	0.06	—
	$\tau_2$	0.69	0.11
	$\tau_3$	0.06	0.013
Stretch-stretch	$F_{a,a}$	1.3	—
	$F_{a,s}$	0.71	—
	$F_{r,1} = F_{d,1}^d$	0.3	0
	$F_{1,1,2}^d$	0.3	0.1
Stretch-bend	$F_{s,s}$	0.15	—
	$F_{a,\alpha} = F_{a,\beta} = F_{r,\alpha} = F_{r,\epsilon} = F_{d,\epsilon}$	0.56	—
	$F_{a,\eta} = F_{r,\xi}$	0.31	—
	$F_{1,\psi} = F_{1,\phi}^d$	0.23	0
	$F_{r,\psi} = F_{d,\phi}$	0.23	0.03
	$F_{1,\theta}^d$	1.13	0
Bend-bend	$F_{\psi_1,\psi_2} = F_{\phi_1,\phi_2}^d$	0.6	−0.03
	$F_{\tau,\psi} = F_{\epsilon,\phi}$	0.3	0
	$F_{\phi_1,\phi_1} = F_{\phi_2,\phi_2}$	0	−0.06

<sup>a</sup> For meaning of symbols, see Fig. 3. Unnumbered indices,  $l$ ,  $\psi$  and  $\phi$  mean  $F_l = F_{l_1} = F_{l_2}$ , etc. <sup>b</sup> In units of mdyn Å<sup>−1</sup> (stretch constants), mdyn rad<sup>−1</sup> (stretch–bend interaction) and mdyn Å rad<sup>−2</sup> (bending and torsion constants). <sup>c</sup> The force constant is the same as for FSA. <sup>d</sup> Interaction between adjacent coordinates.

## NORMAL COORDINATE ANALYSES

We wanted to use force constant calculations as a tool for obtaining more reliable assignments. Because of the similarity in the structure of FSA and SA, it was possible to construct an almost common force field for the two compounds. Except for the force constants directly involving the CF<sub>2</sub> or CH<sub>2</sub> groups, only the C=O stretching force constant is different for FSA and SA. A considerably larger value for this constant had to be used for FSA than for SA in order to reproduce the much higher C=O stretching frequencies for FSA.

We have strictly limited the number of independent force constants, and the final force fields contain 31 parameters given in Table 4. The symbols are defined in Fig. 3. For FSA we have used frequencies from the vapour phase in the calculations whereas

for SA we have used data from solution since no vapour phase data were available for this compound. Hence, the transferability of the force constants are somewhat uncertain since both FSA and SA show considerable frequency shifts when passing from one phase to the other. However, considering the approximations used, the fit between the observed and calculated frequencies is satisfactory.

Only a few normal coordinate analyses are reported for molecules containing CF<sub>2</sub> groups,<sup>14–20</sup> and our final force constants seem to be of reasonable magnitudes compared to them. C–F stretching force constants are previously found from 4.288 for CH<sub>3</sub>CF<sub>2</sub>CH<sub>3</sub><sup>17</sup> to 6.40 mdyn/Å for N=C–CF<sub>2</sub>–CF<sub>2</sub>–C=N<sup>20</sup>, and our value of 5.5 mdyn/Å can therefore be considered as very normal for a

C—F stretching. The CCF bending force constant is almost always found to be *ca.* 1.1 mdyn Å/rad<sup>2</sup> in agreement with the value of 1.07, suitable for FSA. The FCF bending force constant varies from 1.40 (CF<sub>2</sub>Br—CF<sub>2</sub>Br)<sup>14</sup> to 2.331 mdyn Å/rad<sup>2</sup> (CH<sub>3</sub>CF<sub>2</sub>CH<sub>3</sub>),<sup>17</sup> a value of 1.85 thus being very reasonable. Quite different interaction force constants are used for the CF<sub>2</sub> groups.<sup>14–20</sup> Especially the C—C stretching/CCF bending interaction force constant seems to be indeterminable. We have therefore chosen to set this interaction ( $F_{r,\psi}, F_{d,\phi}$ ) constant equal to the CF/CCF interactions ( $F_{1,\psi}, F_{1,\phi}$ ).

The observed and calculated fundamental frequencies are given in Tables 2 (FSA) and 3 (SA) together with the potential energy distribution among the valence coordinates.

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